

## Cesàro Summability of Two-Parameter Trigonometric-Fourier Series\*

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The two-dimensional classical Hardy spaces  $H_p(\mathbf{T} \times \mathbf{T})$  on the bidisc are introduced and it is shown that the maximal operator of the Cesàro means of a distribution is bounded from  $H_p(\mathbf{T} \times \mathbf{T})$  to  $L_p(\mathbf{T}^2)$  ( $3/4 < p \leq \infty$ ) and is of weak type  $(H_1^{\sharp}(\mathbf{T} \times \mathbf{T}), L_1(\mathbf{T}^2))$  where the Hardy space  $H_1^{\sharp}(\mathbf{T} \times \mathbf{T})$  is defined by the hybrid maximal function. As a consequence we obtain that the Cesàro means of a function  $f \in H_1^{\sharp}(\mathbf{T} \times \mathbf{T}) \supset L \log L(\mathbf{T}^2)$  converge a.e. to the function in question. © 1997 Academic Press

### 1. INTRODUCTION

For double trigonometric Fourier series Marcinkiewicz and Zygmund [14] proved that the Cesàro means  $\sigma_{n,m} f$  of a function  $f \in L_1(\mathbf{T}^2)$  converge a.e. to  $f$  as  $n, m \rightarrow \infty$ , provided that the pairs  $(n, m)$  are in a positive cone, i.e., provided that  $2^{-\delta} \leq n/m \leq 2^{\delta}$  for any  $\delta \geq 0$ . A new proof of this result was given by the author [19]. Moreover, Zygmund [24] verified that if  $f \in L \log L(\mathbf{T}^2)$  then the two-parameter Cesàro summability holds.

We proved in [20] and [19] that, in the one-dimensional case, the maximal operator of the Cesàro means of a distribution is bounded from the Hardy–Lorentz space  $H_{p,q}(\mathbf{T})$  to  $L_{p,q}(\mathbf{T})$  if  $3/4 < p \leq \infty$ ,  $0 < q \leq \infty$ , and that, in the two-dimensional case, it is bounded from  $H_{p,q}(\mathbf{T}^2)$  ( $\neq H_{p,q}(\mathbf{T} \times \mathbf{T})$ ) to  $L_{p,q}(\mathbf{T}^2)$  if  $5/6 < p \leq \infty$ ,  $0 < q \leq \infty$ , provided that the supremum in the maximal operator is taken over a positive cone.

In this paper we generalize these results for the unrestricted maximal operator of the two-parameter trigonometric Fourier series. The analogous result for a two-parameter Walsh–Fourier series has been shown by the

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author [21]. The Hardy–Lorentz spaces  $H_{p,q}(\mathbf{T} \times \mathbf{T})$  of distributions are introduced with the  $L_{p,q}(\mathbf{T}^2)$  Lorentz norms of the two-dimensional non-tangential maximal function. Of course,  $H_p(\mathbf{T} \times \mathbf{T}) = H_{p,p}(\mathbf{T} \times \mathbf{T})$  are the usual Hardy spaces ( $0 < p \leq \infty$ ). We will show that the maximal operator of the Cesàro means of a distribution is bounded from  $H_{p,q}(\mathbf{T} \times \mathbf{T})$  to  $L_{p,q}(\mathbf{T}^2)$  ( $3/4 < p \leq \infty$ ,  $0 < q \leq \infty$ ) and is of weak type  $(H_1^\#(\mathbf{T} \times \mathbf{T}), L_1(\mathbf{T}^2))$ , i.e.,

$$\sup_{\gamma > 0} \gamma \lambda(\sup_{n,m \in \mathbf{N}} |\sigma_{n,m} f| > \gamma) \leq C \|f\|_{H_1^\#(\mathbf{T} \times \mathbf{T})} \quad (f \in H_1^\#(\mathbf{T} \times \mathbf{T})).$$

A usual density argument implies then that the Cesàro means  $\sigma_{n,m} f$  converge a.e. to  $f$  as  $n, m \rightarrow \infty$  whenever  $f \in H_1^\#(\mathbf{T} \times \mathbf{T}) \supset L \log L(\mathbf{T}^2)$ . This last result can be found also in Weisz [20].

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## 2. PRELIMINARIES AND NOTATIONS

For a set  $\mathbf{X} \neq \emptyset$  let  $\mathbf{X}^2$  be its Cartesian product  $\mathbf{X} \times \mathbf{X}$  taken with itself, moreover, let  $\mathbf{T} := [-\pi, \pi)$  and  $\lambda$  be the Lebesgue measure. We also use the notation  $|I|$  for the Lebesgue measure of the set  $I$ . We briefly write  $L_p$  instead of the real  $L_p(\mathbf{T}^2, \lambda)$  space while the norm (or quasinorm) of this space is defined by  $\|f\|_p := (\int_{\mathbf{T}^2} |f|^p d\lambda)^{1/p}$  ( $0 < p \leq \infty$ ).

The distribution function of a Lebesgue-measurable function  $f$  is defined by

$$\lambda(\{|f| > \gamma\}) := \lambda(\{x : |f(x)| > \gamma\}) \quad (\gamma \geq 0).$$

The weak  $L_p$  space  $L_p^*(0 < p < \infty)$  consists of all measurable functions  $f$  for which

$$\|f\|_{L_p^*} := \sup_{\gamma > 0} \gamma \lambda(\{|f| > \gamma\})^{1/p} < \infty$$

while we set  $L_\infty^* = L_\infty$ .

The spaces  $L_p^*$  are special cases of the more general Lorentz spaces  $L_{p,q}$ . In their definition another concept is used. For a measurable function  $f$  the *non-increasing rearrangement* is defined by

$$\tilde{f}(t) := \inf\{\gamma : \lambda(\{|f| > \gamma\}) \leq t\}.$$

Lorentz space  $L_{p,q}$  is defined as follows: for  $0 < p < \infty$ ,  $0 < q < \infty$

$$\|f\|_{p,q} := \left( \int_0^\infty \tilde{f}(t)^q t^{q/p} \frac{dt}{t} \right)^{1/q}$$

while for  $0 < p \leq \infty$

$$\|f\|_{p, \infty} := \sup_{t > 0} t^{1/p} \tilde{f}(t).$$

Let

$$L_{p, q} := L_{p, q}(\mathbf{T}^j, \lambda) := \{f : \|f\|_{p, q} < \infty\} \quad (j = 1, 2).$$

One can show the following equalities:

$$L_{p, p} = L_p, \quad L_{p, \infty} = L_p^* \quad (0 < p \leq \infty)$$

(see e.g. Bennett, Sharpley [1] or Bergh, Löfström [2]).

Let  $f$  be a distribution on  $C^\infty(\mathbf{T}^2)$  (briefly  $f \in \mathcal{D}'(\mathbf{T}^2) = \mathcal{D}'$ ). The  $(n, m)$ th Fourier coefficient is defined by  $\hat{f}(n, m) := f(e^{-inx} e^{-imy})$  where  $i = \sqrt{-1}$ . In special case, if  $f$  is an integrable function then

$$\hat{f}(n, m) = \frac{1}{(2\pi)^2} \int_{\mathbf{T}} \int_{\mathbf{T}} f(x, y) e^{-inx} e^{-imy} dx dy.$$

For simplicity, we assume that, for a distribution  $f \in \mathcal{D}'$ , we have  $\hat{f}(n, 0) = \hat{f}(0, n) = 0$  ( $n \in \mathbf{N}$ ). Denote by  $s_{n, m} f$  the  $(n, m)$ th partial sum of the Fourier series of a distribution  $f$ , namely,

$$s_{n, m} f(x) := \sum_{k=-n}^n \sum_{l=-m}^m \hat{f}(k, l) e^{ikx} e^{ily}.$$

For  $f \in \mathcal{D}'$  and  $z_1 := re^{ix}$ ,  $z_2 := se^{iy}$  ( $0 < r, s < 1$ ) let

$$u(z_1, z_2) = u(re^{ix}, se^{iy}) := (f * P_r \times P_s)(x, y)$$

where  $*$  denotes the convolution and

$$P_r(x) := \sum_{k=-\infty}^{\infty} r^{|k|} e^{ikx} = \frac{1-r^2}{1+r^2-2r \cos x} \quad (x \in \mathbf{T})$$

is the Poisson kernel. It is easy to show that  $u(z_1, z_2)$  is a biharmonic function on the bidisc and

$$u(re^{ix}, se^{iy}) = \sum_{k=-\infty}^{\infty} \sum_{l=-\infty}^{\infty} \hat{f}(k, l) r^{|k|} s^{|l|} e^{ikx} e^{ily}$$

with absolute and uniform convergence (see e.g. Gundy, Stein [12], Edwards [8]).

Let  $0 < \alpha < 1$  be an arbitrary number. We denote by  $\Omega_\alpha(x)$  ( $x \in \mathbf{T}$ ) the region bounded by two tangents to the circle  $|z| = \alpha$  from  $e^{ix}$  and the longer arc of the circle included between the points of tangency. The non-tangential maximal function is defined by

$$u_{\alpha, \beta}^*(x, y) := \sup_{z_1 \in \Omega_\alpha(x)} \sup_{z_2 \in \Omega_\beta(y)} |u(z_1, z_2)| \quad (0 < \alpha, \beta < 1).$$

For  $0 < p, q \leq \infty$  the *Hardy–Lorentz space*  $H_{p, q}(\mathbf{T} \times \mathbf{T}) = H_{p, q}$  consists of all distributions  $f$  for which  $u_{\alpha, \beta}^* \in L_{p, q}$  and set

$$\|f\|_{H_{p, q}} := \|u_{1/2, 1/2}^*\|_{p, q}.$$

It is known that if  $f \in H_p$  ( $0 < p < \infty$ ) then  $f(x, y) = \lim_{r, s \rightarrow 1} u(re^{ix}, se^{iy})$  in the sense of distributions (see Gundy, Stein [12]).

Let us introduce the hybrid Hardy spaces. For  $f \in L_1(\mathbf{T}^2)$  and  $z := re^{ix}$  ( $0 < r < 1$ ) let

$$v(z, y) = v(re^{ix}, y) := \frac{1}{2\pi} \int_{\mathbf{T}} f(t, y) P_r(x - t) dt$$

and

$$v_\alpha^+(x, y) := \sup_{z \in \Omega_\alpha(x)} |v(z, y)| \quad (0 < \alpha < 1).$$

We say that  $f \in L_1(\mathbf{T}^2)$  is in the hybrid Hardy–Lorentz space  $H_{p, q}^\#(\mathbf{T} \times \mathbf{T}) = H_{p, q}^\#$  if

$$\|f\|_{H_{p, q}^\#} := \|v_{1/2}^+\|_{p, q} < \infty.$$

The equivalences  $\|u_{\alpha, \beta}^*\|_{p, q} \sim \|u_{1/2, 1/2}^*\|_{p, q}$ ,  $\|v_\alpha^+\|_{p, q} \sim \|v_{1/2}^+\|_{p, q}$  ( $0 < p, q < \infty$ ,  $0 < \alpha, \beta < 1$ ) and  $H_{p, q} \sim H_{p, q}^\# \sim L_{p, q}$  ( $1 < p < \infty$ ,  $0 < q \leq \infty$ ) were proved in Fefferman, Stein [9], Gundy, Stein [12] and Lin [13]. Note that in case  $p = q$  the usual definition of Hardy spaces  $H_{p, p} = H_p$  and  $H_{p, p}^\# = H_p^\#$  are obtained. For other equivalent definitions we call for Gundy, Stein [12], Gundy [11] and Chang, Fefferman [4].

In this paper the constants  $C$  are absolute constants and the constants  $C_p$  (resp.  $C_{p, q}$ ) are depending only on  $p$  (resp.  $p$  and  $q$ ) and may denote different constants in different contexts.

Recall that, in the one-dimensional case,  $L_1(\mathbf{T}) \subset H_{1, \infty}(\mathbf{T})$  and  $L \log L(\mathbf{T}) \subset H_1(\mathbf{T})$ , more exactly,

$$\|f\|_{H_{1, \infty}(\mathbf{T})} = \sup_{\gamma > 0} \gamma \lambda(u_{1/2}^* > \gamma) \leq C \|f\|_1 \quad (f \in L_1(\mathbf{T})) \quad (1)$$

and

$$\|f\|_{H_1(\mathbf{T})} \leq C + C \| |f| \log^+ |f| \|_1 \quad (f \in L \log L(\mathbf{T})) \quad (2)$$

where  $\log^+ u = 1_{\{u > 1\}} \log u$  (see Fefferman, Stein [9] and Stein [18]).

These results are generalized for two parameters in the following way.

**THEOREM 1.** *We have  $L \log L \subset H_1^\sharp \subset H_{1, \infty}$  more exactly,*

$$\|f\|_{H_{1, \infty}} = \sup_{\gamma > 0} \gamma \lambda(u_{1/2, 1/2}^* > \gamma) \leq C \|f\|_{H_1^\sharp} \quad (f \in H_1^\sharp) \quad (3)$$

and

$$\|f\|_{H_1^\sharp} \leq C + C \| |f| \log^+ |f| \|_1 \quad (f \in L \log L). \quad (4)$$

*Proof.* Applying Fubini's theorem, (1) and the positivity of the Poisson kernel we have

$$\begin{aligned} & \lambda \left( (x, y) : \sup_{re^{iw} \in \Omega_{1/2}(x)} \sup_{se^{iw} \in \Omega_{1/2}(y)} \left| \int_{\mathbf{T}} \int_{\mathbf{T}} f(t, u) P_r(v-t) P_s(w-u) dt du \right| > \gamma \right) \\ & \leq \lambda \left( (x, y) : \sup_{se^{iw} \in \Omega_{1/2}(y)} \int_{\mathbf{T}} \left( \sup_{re^{iw} \in \Omega_{1/2}(x)} \left| \int_{\mathbf{T}} f(t, u) P_r(v-t) dt \right| \right) \right. \\ & \quad \left. \times P_s(w-u) du > \gamma \right) \\ & = \int_{\mathbf{T}} \int_{\mathbf{T}} 1_{\{\sup_{se^{iw} \in \Omega_{1/2}(\cdot)} \int_{\mathbf{T}} (\sup_{re^{iw} \in \Omega_{1/2}(\cdot)} |\int_{\mathbf{T}} f(t, u) P_r(v-t) dt|) P_s(w-u) du > \gamma\}}(x, y) dy dx \\ & \leq \frac{C}{\gamma} \int_{\mathbf{T}} \int_{\mathbf{T}} \sup_{re^{iw} \in \Omega_{1/2}(x)} \left| \int_{\mathbf{T}} f(t, y) P_r(v-t) dt \right| dy dx \\ & = \frac{C}{\gamma} \int_{\mathbf{T}} \int_{\mathbf{T}} v_{1/2}^+(x, y) dx dy \end{aligned}$$

which proves (3). (4) comes easily from (2). ■

### 3. ATOMIC DECOMPOSITION AND BOUNDED OPERATORS ON HARDY SPACES

A *generalized interval* on  $\mathbf{T}$  is either an interval  $I \subset \mathbf{T}$  or  $I = [-\pi, x) \cup [y, \pi)$ . A *generalized rectangle* on  $\mathbf{T}^2$  is the Descartes product  $I \times J$  of two generalized intervals.

A function  $a \in L_2$  is a  $p$ -atom if

(i)  $\text{supp } a \subset F$  for an open set  $F \subset \mathbf{T}^2$

(ii)  $\|a\|_2 \leq \lambda(F)^{1/2-1/p}$

(iii)  $a = \sum_R \lambda_R a_R$  where the  $\lambda_R$ 's are real numbers and the  $a_R$ 's are functions (called "elementary particles") satisfying

( $\alpha$ )  $\text{supp } a_R \subset R$  for any generalized rectangle  $R = I \times J \subset F$

( $\beta$ )

$$\left\| \frac{\partial^N a_R(x, y)}{\partial x^N} \right\|_\infty \leq \frac{C}{\sqrt{|R|} |I|^N} \quad \text{and} \quad \left\| \frac{\partial^N a_R(x, y)}{\partial y^N} \right\|_\infty \leq \frac{C}{\sqrt{|R|} |J|^N}$$

for all  $N \leq [2/p - 1/2]$

( $\gamma$ ) for all  $x, y \in \mathbf{T}$  and all  $M \leq [2/p - 3/2]$

$$\int_{\mathbf{T}} a_R(x, y) x^M dx = \int_{\mathbf{T}} a_R(x, y) y^M dy = 0$$

( $\delta$ )  $\left( \sum_R \lambda_R^2 \right)^{1/2} \leq \lambda(F)^{1/2-1/p}$ .

If  $a \in L_2$  satisfies (i) with a generalized rectangle  $F$ , (ii) and ( $\gamma$ ) then  $a$  is called a *rectangle  $p$ -atom*.

The basic result of the atomic decomposition is stated as follows (see Chang, Fefferman [4], Fefferman [10], Wilson [23] and also Weisz [21]).

**THEOREM A.** *A distribution  $f$  is in  $H_p$  ( $0 < p \leq 1$ ) if and only if there exist a sequence  $(a_k, k \in \mathbf{N})$  of  $p$ -atoms and a sequence  $(\mu_k, k \in \mathbf{N})$  of real numbers such that*

$$\sum_{k=0}^{\infty} \mu_k a_k = f \quad \text{in the sense of distributions,} \tag{5}$$

$$\sum_{k=0}^{\infty} |\mu_k|^p < \infty.$$

Moreover, the following equivalence of norms holds:

$$\|f\|_{H_p} \sim \inf \left( \sum_{k=0}^{\infty} |\mu_k|^p \right)^{1/p} \quad (6)$$

where the infimum is taken over all decompositions of  $f$  of the form (5).

If  $I$  is a generalized interval then let  $2^r I$  be the generalized interval with the same center as  $I$  and with length  $2^r |I|$  ( $r \in \mathbf{N}$ ). For a generalized rectangle  $R = I \times J$  let  $2^r R = 2^r I \times 2^r J$ .

Using Theorem A, (3) and the interpolation results given in Lin [13] we can prove the following theorem similarly to the Theorem of Fefferman [10] (see also Corollary 1 in Weisz [21]).

**THEOREM 2.** *Suppose that the operator  $T$  is sublinear and  $p_0 < p \leq 1$ . Furthermore, assume that there exists  $\delta > 0$  such that for every rectangle  $p$ -atom  $a$  supported on the generalized rectangle  $R$  and for every  $r \geq 2$  one has*

$$\int_{\mathbf{T}^2 \setminus 2^r R} |Ta|^p d\lambda \leq C_p 2^{-\delta r} \quad (7)$$

where  $C_p$  is a constant depending only on  $p$ . If  $T$  is bounded from  $L_p$  to  $L_p$  ( $p = 2, \infty$ ) then

$$\|Tf\|_{p,q} \leq C_{p,q} \|f\|_{H,q} \quad (f \in H_{p,q})$$

for every  $p_0 < p < \infty$  and  $0 < q \leq \infty$ . Specially,  $T$  is of weak type  $(H_1^\#, L_1)$ , i.e. if  $f \in H_1^\#$  then

$$\|Tf\|_{1,\infty} = \sup_{\gamma \geq 0} \gamma \lambda(|Tf| > \gamma) \leq C \|f\|_{H_{1,\infty}} \leq C \|f\|_{H_1^\#}.$$

#### 4. CESÀRO SUMMABILITY OF TWO-PARAMETER TRIGONOMETRIC-FOURIER SERIES

For  $n, m \in \mathbf{N}$  and a distribution  $f$  the Cesàro mean of order  $(n, m)$  of the Fourier series of  $f$  is given by

$$\sigma_{n,m} f := \frac{1}{n+1} \frac{1}{m+1} \sum_{k=0}^n \sum_{l=0}^m s_{k,l} f = f * (K_n \times K_m) \quad (n, m \in \mathbf{N})$$

where  $K_n$  is the Fejér kernel of order  $n$ . It is shown in Zygmund [24] that

$$0 \leq K_n(t) \leq \frac{\pi^2}{(n+1)t^2} \quad (0 < |t| < \pi) \tag{8}$$

and

$$\int_{\mathbf{T}} K_n(t) dt = \pi. \tag{9}$$

For a distribution  $f$  we consider the maximal operator of the Cesàro means

$$\sigma^* f := \sup_{n, m \in \mathbf{N}} |\sigma_{n, m} f|$$

and prove our following main result.

**THEOREM 3.** *There are absolute constants  $C$  and  $C_{p, q}$  such that*

$$\|\sigma^* f\|_{p, q} \leq C_{p, q} \|f\|_{H_{p, q}} \quad (f \in H_{p, q}) \tag{10}$$

for every  $3/4 < p < \infty$  and  $0 < q \leq \infty$ . Especially, if  $f \in H_1^\#$  then

$$\lambda(\sigma^* f > \gamma) \leq \frac{C}{\gamma} \|f\|_{H_1^\#} \quad (\gamma > 0). \tag{11}$$

*Proof.* It is proved in Zygmund [24] that

$$\|\sigma^* f\|_p \leq C_p \|f\|_p \quad (1 < p \leq \infty). \tag{12}$$

So, by Theorem 2, the proof of Theorem 3 will be complete if we show that the operator  $\sigma^*$  satisfies (7) for each  $3/4 < p \leq 1$ .

Let  $a$  be an arbitrary rectangle  $p$ -atom with support  $R = I \times J$  and

$$2^{-K-1} < |I|/\pi \leq 2^{-K}, \quad 2^{-L-1} < |J|/\pi \leq 2^{-L} \quad (K, L \in \mathbf{N}).$$

We can suppose that the center of  $R$  is zero. In this case

$$[-\pi 2^{-K-2}, \pi 2^{-K-2}] \subset I \subset [-\pi 2^{-K-1}, \pi 2^{-K-1}]$$

and

$$[-\pi 2^{-L-2}, \pi 2^{-L-2}] \subset J \subset [-\pi 2^{-L-1}, \pi 2^{-L-1}].$$



To prove (7) for the operator  $\sigma^*$  we have to integrate  $|\sigma^*a|^p$  over

$$\begin{aligned} \mathbf{T}^2 \setminus 2^r R &= (\mathbf{T} \setminus 2^r I) \times J \cup (\mathbf{T} \setminus 2^r I) \times (\mathbf{T} \setminus J) \\ &\cup I \times (\mathbf{T} \setminus 2^r J) \cup (\mathbf{T} \setminus I) (\mathbf{T} \setminus 2^r J) \end{aligned}$$

where  $r \geq 2$  is an arbitrary integer. We do this in four steps.

*Step 1: Integrating over  $(\mathbf{T} \setminus 2^r I) \times J$ .* Obviously,

$$\begin{aligned} &\int_{\mathbf{T} \setminus 2^r I} \int_J |\sigma^*a(x, y)|^p dx dy \\ &\leq \sum_{|i|=2^{r-2}}^{2^K-1} \int_{\pi i 2^{-K}}^{\pi(i+1)2^{-K}} \int_J |\sigma^*a(x, y)|^p dx dy \\ &\leq \sum_{|i|=2^{r-2}}^{2^K-1} \int_{\pi i 2^{-K}}^{\pi(i+1)2^{-K}} \int_J \sup_{n \geq r_i, m \in \mathbf{N}} |\sigma_{n,m}a(x, y)|^p dx dy \\ &\quad + \sum_{|i|=2^{r-2}}^{2^K-1} \int_{\pi i 2^{-K}}^{\pi(i+1)2^{-K}} \int_J \sup_{n < r_i, m \in \mathbf{N}} |\sigma_{n,m}a(x, y)|^p dx dy \\ &= (A) + (B) \end{aligned} \tag{13}$$

where  $r_i := \lceil 2^K/i^\alpha \rceil$  ( $i \in \mathbf{N}$ ) with  $\alpha > 0$  chosen later.

It is easy to see that

$$\frac{1}{(x-t)^2} \leq \frac{1}{(\pi i 2^{-K} - \pi 2^{-K-1})^2} \leq \frac{4}{\pi^2} \frac{2^{2K}}{i^2} \tag{14}$$

if  $x \in [\pi i 2^{-K}, \pi(i+1)2^{-K})$  ( $|i| \geq 1$ ) and  $t \in I$ . Hence, by (8),

$$\begin{aligned} |\sigma_{n,m}a(x, y)| &= \left| \int_I \int_J a(t, u) K_n(x-t) K_m(y-u) dt du \right| \\ &\leq \int_I \left| \int_J a(t, u) K_m(y-u) du \right| K_n(x-t) dt \\ &\leq C \int_I \left| \int_J a(t, u) K_m(y-u) du \right| \frac{1}{(n+1)(x-t)^2} dt \\ &\leq \frac{C 2^{2K}}{(n+1)i^2} \int_I \left| \int_J a(t, u) K_m(y-u) du \right| dt. \end{aligned} \tag{15}$$

By Hölder's inequality,

$$\begin{aligned} & \int_J \sup_{n \geq r_i, m \in \mathbf{N}} |\sigma_{n,m} a(x, y)|^p dy \\ & \leq \frac{C_p 2^{2Kp}}{(r_i + 1)^p i^{2p}} |J|^{1-p} \left( \int_I \int_J \sup_{m \in \mathbf{N}} \left| \int_J a(t, u) K_m(y-u) du \right| dy dt \right)^p. \end{aligned} \quad (16)$$

Using again Hölder's inequality and (12) for one dimension and for a fixed  $t$ , we obtain

$$\begin{aligned} & \int_J \sup_{m \in \mathbf{N}} \left| \int_J a(t, u) K_m(y-u) du \right| dy \\ & \leq |J|^{1/2} \left( \int_{\mathbf{T}} \sup_{m \in \mathbf{N}} \left| \int_J a(t, u) K_m(y-u) du \right|^2 dy \right)^{1/2} \\ & \leq C |J|^{1/2} \left( \int_J |a(t, y)|^2 dy \right)^{1/2}. \end{aligned}$$

By the definition of the rectangle  $p$ -atom we conclude that

$$\begin{aligned} & \int_I \int_J \sup_{m \in \mathbf{N}} \left| \int_J a(t, u) K_m(y-u) du \right| dy dt \\ & \leq C |J|^{1/2} |I|^{1/2} \left( \int_{\mathbf{T}} \int_{\mathbf{T}} |a(t, y)|^2 dy dt \right)^{1/2} \\ & \leq C_p 2^{-K+K/p-L+L/p}. \end{aligned} \quad (17)$$

Using the value of  $r_i$  we can establish that

$$\begin{aligned} (A) & \leq C_p \sum_{i=2^{r-2}}^{2^K-1} 2^{-K} \frac{2^{2Kp}}{(r_i + 1)^p i^{2p}} 2^{-L+Lp} 2^{-Kp+K-Lp+L} \\ & \leq C_p \sum_{i=2^{r-2}}^{2^K-1} \frac{1}{i^{2p-\alpha p}} \leq C_{p,\alpha} 2^{-r(2p-\alpha p-1)} \end{aligned} \quad (18)$$

provided that

$$\alpha < \frac{2p-1}{p} (\leq 1). \quad (19)$$

Now let us consider (B). It is known that

$$\begin{aligned}\sigma_{n,m}a(x,y) &= \int_I \left( \int_J a(t,u) K_m(y-u) du \right) K_n(x-t) dt \\ &= \frac{1}{2\pi} \sum_{|k|=0}^n \left( 1 - \frac{|k|}{n+1} \right) \int_I \left( \int_J a(t,u) K_m(y-u) du \right) e^{ikt} dt e^{ikx}.\end{aligned}$$

By the definition of the atom,

$$\begin{aligned}& \left| \int_I \left( \int_J a(t,u) K_m(y-u) du \right) e^{ikt} dt \right| \\ & \leq \left| \int_I \left( \int_J a(t,u) K_m(y-u) du \right) (e^{ikt} - 1) dt \right| \\ & \leq \int_I \left| \int_J a(t,u) K_m(y-u) du \right| |k| |t| dt \\ & \leq |I| |k| \int_I \left| \int_J a(t,u) K_m(y-u) du \right| dt.\end{aligned}$$

Therefore

$$\begin{aligned}& \sup_{n < r_i, m \in \mathbf{N}} |\sigma_{n,m}a(x,y)| \\ & \leq C \sup_{n < r_i, m \in \mathbf{N}} \sum_{|k|=0}^n \frac{n+1-|k|}{n+1} |I| |k| \int_I \left| \int_J a(t,u) K_m(y-u) du \right| dt \\ & \leq C_p \sum_{k=0}^{r_i} (r_i - k) 2^{-K} \int_I \sup_{m \in \mathbf{N}} \left| \int_J a(t,u) K_m(y-u) du \right| dt \\ & \leq C_p r_i^2 2^{-K} \int_I \sup_{m \in \mathbf{N}} \left| \int_J a(t,u) K_m(y-u) du \right| dt.\end{aligned}\tag{20}$$

It follows similarly to (16) and (17) that

$$\begin{aligned}(B) & \leq C_p \sum_{i=2^{r-2}}^{2^K-1} 2^{-Kr} i^{2p} 2^{-Kp} |J|^{1-p} \\ & \quad \times \left( \int_I \int_J \sup_{m \in \mathbf{N}} \left| \int_J a(t,u) K_m(y-u) du \right| dy dt \right)^p \\ & \leq C_p \sum_{i=2^{r-2}}^{2^K-1} 2^{-K} 2^{2Kp} i^{-2\alpha p} 2^{-Kp} 2^{-L+Lp} 2^{-Kp+K-Lp+L} \\ & \leq C_p \sum_{i=2^{r-2}}^{2^K-1} i^{-2\alpha p} \leq C_{p,\alpha} 2^{-r(2\alpha p-1)}\end{aligned}\tag{21}$$

whenever

$$\alpha > \frac{1}{2p}. \quad (22)$$

The number  $\alpha$  satisfies (19) and (22) if and only if  $3/4 < p \leq 1$ .

Combining (18) and (21) we can establish that, for  $3/4 < p \leq 1$ ,

$$\int_{\mathbf{T} \setminus 2^r I} \int_J |\sigma^* a(x, y)|^p dx dy \leq C_p 2^{-\delta r} \quad (23)$$

where  $C_p$  depends only on  $p$ .

*Step 2. Integrating over  $(\mathbf{T} \setminus 2^r I) \times (\mathbf{T} \setminus J)$ .* Similarly to (13),

$$\begin{aligned} & \int_{\mathbf{T} \setminus 2^r I} \int_{\mathbf{T} \setminus J} |\sigma^* a(x, y)|^p dx dy \\ & \leq \sum_{|i|=2^{r-2}}^{2^K-1} \sum_{|j|=1}^{2^L-1} \int_{\pi i 2^{-K}}^{\pi(i+1)2^{-K}} \int_{\pi j 2^{-L}}^{\pi(j+1)2^{-L}} \sup_{n \geq r_i, m \geq s_j} |\sigma_{n,m} a(x, y)|^p dx dy \\ & \quad + \sum_{|i|=2^{r-2}}^{2^K-1} \sum_{|j|=1}^{2^L-1} \int_{\pi i 2^{-K}}^{\pi(i+1)2^{-K}} \int_{\pi j 2^{-L}}^{\pi(j+1)2^{-L}} \sup_{n < r_i, m \geq s_j} |\sigma_{n,m} a(x, y)|^p dx dy \\ & \quad + \sum_{|i|=2^{r-2}}^{2^K-1} \sum_{|j|=1}^{2^L-1} \int_{\pi i 2^{-K}}^{\pi(i+1)2^{-K}} \int_{\pi j 2^{-L}}^{\pi(j+1)2^{-L}} \sup_{n \geq r_i, m < s_j} |\sigma_{n,m} a(x, y)|^p dx dy \\ & \quad + \sum_{|i|=2^{r-2}}^{2^K-1} \sum_{|j|=1}^{2^L-1} \int_{\pi i 2^{-K}}^{\pi(i+1)2^{-K}} \int_{\pi j 2^{-L}}^{\pi(j+1)2^{-L}} \sup_{n < r_i, m < s_j} |\sigma_{n,m} a(x, y)|^p dx dy \\ & = (C) + (D) + (E) + (F) \end{aligned}$$

where  $r_i := [2^K/i^\alpha]$  and  $s_j := [2^L/j^\alpha]$  ( $i, j \in \mathbf{N}$ ) with  $\alpha > 0$  chosen later.

Similarly to (15) and (17),

$$\begin{aligned} |\sigma_{n,m} a(x, y)| & = \left| \int_I \int_J a(t, u) K_n(x-t) K_m(y-u) dt du \right| \\ & \leq C \int_I \int_J |a(t, u)| \frac{1}{(n+1)(x-t)^2} \frac{1}{(m+1)(y-u)^2} dt du \\ & \leq \frac{C 2^{2K} 2^{2L}}{(n+1)(m+1) i^2 j^2} |I|^{1/2} |J|^{1/2} \left( \int_I \int_J |a(t, u)|^2 dt du \right)^{1/2} \\ & \leq \frac{C_p 2^{K+L+K/p+L/p}}{(n+1)(m+1) i^2 j^2}. \end{aligned}$$

Consequently,

$$\begin{aligned}
 (C) &\leq C_p \sum_{i=2^{r-2}}^{2^K-1} \sum_{j=1}^{2^L-1} 2^{-K-L} \frac{2^{Kp+Lp+K+L}}{(r_i+1)^p (s_j+1)^p i^{2p} j^{2p}} \\
 &\leq C_p \sum_{|i|=2^{r-2}}^{2^K-1} \sum_{|j|=1}^{2^L-1} \frac{1}{i^{2p-\alpha p} j^{2p-\alpha p}} \leq C_{p,\alpha} 2^{-r(2p-\alpha p-1)} \quad (24)
 \end{aligned}$$

whenever (19) holds.

We get from (14), (17) and (20) that

$$\begin{aligned}
 \sup_{n < r_i, m \geq s_j} |\sigma_{n,m} a(x, y)| &\leq C_p r_i^2 2^{-K} \sup_{m \geq s_j} \int_I \int_J |a(t, u)| \frac{1}{(m+1)(y-u)^2} dt du \\
 &\leq C_p r_i^2 2^{-K} \frac{1}{(s_j+1)} \frac{2^{2L}}{j^2} 2^{-K-L+K/p+L/p} \\
 &\leq C_p \frac{2^{K/p+L/p}}{i^{2\alpha} j^{2-\alpha}}.
 \end{aligned}$$

Hence

$$(D) \leq C_p \sum_{i=2^{r-2}}^{2^K-1} \sum_{j=1}^{2^L-1} 2^{-K-L} \frac{2^{K+L}}{i^{2\alpha p} j^{(2-\alpha)p}} \leq C_{p,\alpha} 2^{-r(2\alpha p-1)} \quad (25)$$

if (19) and (22) are satisfied.

Similarly,

$$\sup_{n \geq r_i, m < s_j} |\sigma_{n,m} a(x, y)| \leq C_p \frac{2^{K/p+L/p}}{j^{2\alpha} i^{2-\alpha}}$$

and

$$(E) \leq C_p \sum_{i=2^{r-2}}^{2^K-1} \sum_{j=1}^{2^L-1} 2^{-K-L} \frac{2^{K+L}}{j^{2\alpha p} i^{(2-\alpha)p}} \leq C_{p,\alpha} 2^{-r(2p-\alpha p-1)} \quad (26)$$

provided that (19) and (22) are true.

We know that

$$\sigma_{n,m} a(x, y) = \sum_{|k|=0}^n \sum_{|l|=0}^m \left(1 - \frac{|k|}{n+1}\right) \left(1 - \frac{|l|}{m+1}\right) \hat{a}(k, l) e^{ikx +ily}.$$

Next, we establish that

$$\begin{aligned}
|\hat{a}(k, l)| &= \left| \frac{1}{(2\pi)^2} \int_I \int_J a(x, y)(e^{-ikx} - 1)(e^{-ily} - 1) dx dy \right| \\
&\leq \frac{1}{(2\pi)^2} \int_I \int_J |a(x, y)| |kx| |ly| dx dy \\
&\leq \frac{1}{(2\pi)^2} |k| |l| |I| |J| |I|^{1/2} |J|^{1/2} \left( \int_I \int_J |a(x, y)|^2 dx dy \right)^{1/2} \\
&\leq C_p |k| |l| 2^{-2K-2L+K/p+L/p}.
\end{aligned}$$

So

$$\begin{aligned}
\sup_{n < r_i, m < s_j} |\sigma_{n, m} a(x, y)| &\leq C \sup_{n < r_i, m < s_j} \sum_{|k|=0}^n \sum_{|l|=0}^m \frac{n+1-|k|}{n+1} \\
&\quad \times \frac{m+1-|l|}{m+1} |\hat{a}(k, l)| \\
&\leq C_p \sum_{k=0}^{r_i} \sum_{l=0}^{s_j} (r_i - k)(s_j - l) 2^{-2K-2L+K/p+L/p} \\
&\leq C_p r_i^2 s_j^2 2^{-2K-2L+K/p+L/p}.
\end{aligned}$$

A simple calculation shows that

$$\begin{aligned}
(F) &\leq C_p \sum_{i=2^{r-2}}^{2^K-1} \sum_{j=1}^{2^L-1} 2^{-K-L} r_i^{2p} s_j^{2p} 2^{-2Kp-2Lp+K+L} \\
&\leq C_p \sum_{i=2^{r-2}}^{2^K-1} \sum_{j=1}^{2^L-1} \frac{1}{i^{2\alpha p} j^{2\alpha p}} \leq C_p \alpha 2^{-r(2\alpha p-1)}
\end{aligned}$$

supposed that (19) holds.

Combining (24)–(27) we can see that, for  $3/4 < p \leq 1$ ,

$$\int_{\mathbf{T} \setminus 2^r I} \int_{\mathbf{T} \setminus J} |\sigma^* a(x, y)|^p dx dy \leq C_p 2^{-\delta r} \quad (28)$$

where  $C_p$  depends only on  $p$ .

Steps 3 and 4, the integration over  $I \times (\mathbf{T} \setminus 2^r J)$  and over  $(\mathbf{T} \setminus I) \times (\mathbf{T} \setminus 2^r J)$ , are analogous to Steps 1 and 2.

Taking into account (23) and (28), we have proved (7) and also the theorem. ■

Note that Theorem 3 was proved by the author for the one-parameter case, for the restricted maximal operator and  $H_p(\mathbf{T}^2)$  and, moreover, for the two-parameter Walsh–Fourier series (see [20], [19], [21]).

We suspect that Theorem 3 for  $p \leq 3/4$  is not true though we could not find any counterexample.

It is easy to show that the two-dimensional trigonometric polynomials are dense in  $H_1^\#$ . Hence (11) and the usual density argument (see Marcinkiewicz, Zygmund [14]) imply

COROLLARY 1. *If  $f \in H_1^\#$  then*

$$\sigma_{n,m} f \rightarrow f \quad \text{a.e. as } \min(n, m) \rightarrow \infty.$$

Note that this corollary is proved in [20] with another method.

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